APPENDCES C
PUBLICATIONS

## First Author

**S. Ruamruk**, K. Singsoog, P. Pilasuta, S. Paengson, W. Namhongsa, M. Rittiruam and T. Seetawan. (2018). Electronic structure and thermoelectric properties of  $Bi_{2-}$   $_xSb_xTe_3$ 

(x = 0, 0.33, 0.67) by first principle calculation. Materials Today: Proceedings 5 (2018) 14150–14154

and the section of th



Available online at www.sciencedirect.com

# ScienceDirect





www.materialstoday.com/proceedings

## **SACT 2016**

# Electronic structure and thermoelectric properties of Bi<sub>2-x</sub>Sb<sub>x</sub>Te<sub>3</sub> (x = 0, 0.33, 0.67) by first principle calculation

S. Ruamruk $^{a,b}$ , K. Singsoog $^{a,b}$ , P. Pilasuta $^{a,b}$ , S. Paengson $^{a,b}$ , W. Namhongsa $^{a,b}$ , M. Rittiruam $^{b,c}$ , T. Seetawan $^{a,b,c,*}$ 

<sup>a</sup>Thermoelectirc Research Laboratory, Center of Excellence on Alternative Energy, Research and Development Institution,

Inermoelective Research Laboratory, Center of Excellence on Alternative Energy, Research and Development Institution,
Sakon Nakhon Rajabhat University 4700, Thailand

bProgram of Physics, Faculty of Science and Technology, Sakon Nakhon Rajabhat University, Sakon Nakhon 1700, Thailand

cSimulation Research Laboratory, Center of Excellence on Alternative Energy, Research and Development Institution,
Sakon Nakhon Rajabhat University, Sakon Nakhon 47000, Thailand

#### Abstract

The electronic structure and thermoelectric properties of  $\mathrm{Bi}_{2x}\mathrm{Sb}_x\mathrm{Te}_3$  (x = 0, 0.33, 0.67) were systematically investigated by the first-principles calculations and Boltzmann transport theory based on QUANTUM ERPREESO and BoltzTraP. According to the calculated electrical conductivity, Seebeck coefficient and thermal conductivity with respected to empirical scattering by substituted Bi by Sb for rhombohedral structure exhibit diverse anisotropic thermoelectric properties. It is found that the  $Bi_{2,x}Sb_xTe_3$  (x = 0, 0.33, 0.67) show n- and p-type substituting better thermoelectric performance compared to  $Bi_2Te_3$ . We would like to be beneficial to stimulate further theoretical and experimental works.

© 2017 Elsevier Ltd. All rights reserved.

Selection and/or Peer-review under responsibility of SACT 2016.

Keywords: Thermoelectric properties; Boltzmann transport theory; Bi2xSbxTe3; density functional theory; semiconductor

## 1. Introduction

Thermoelectric (TE) materials have been earned heat into electricity as well as it much attention for their potential application in power generate and refrigerate [1]. The quality of TE materials is related to a parameter called the dimensionless figure of merit:  $ZT = S^2 \sigma T / (\kappa_e + \kappa_L)$ , where S is the Seebeck coefficient, T is absolute

<sup>\*</sup> Corresponding author. Tel.: +6642-744-319; fax: +6642-744-319. E-mail address: t\_seetawan@snru.ac.th

temperature,  $\sigma$  is electrical conductivity,  $\kappa_e$  is the thermal conductivity and  $\kappa_t$  is the lattice thermal conductivity [2]. Bismuth Telluride, (Bi<sub>2</sub>Te<sub>3</sub>) has been significant role thermoelectric technology due to its low Seebeck coefficient, high electrical conductivity and low thermal conductivity [3]. We proposed to improve performance of this material by substitution Antimony in Bi-site. In addition, the TE properties of TE materials can be investigated by the Boltzmann transport theory [4, 5]. In this work, we studies. Theoretical substituted Bi by Sb to investigate the electronic structure and enhanced TE properties of Bi<sub>2</sub>Te<sub>3</sub> material by using density functional theory and BoltzTraP simulation.

#### 2. Computational details

The  $Bi_{2-x}Sb_xTe_3$  (x = 0, 0.33, 0.67) cluster models were designed by using a space group number 166, lattice parameters a = 4.38 Å, b = 4.38 Å and c = 30.49 Å, as show in Fig. 1. The electronic structure was calculated by the density functional theory based on QUANTUM ESPRESSO [6]. The exchange correlation function, energy convergence limit set as 10<sup>-8</sup> Ry, and energy cutoff 40 Ry are performed [7]. The density of states was calculated using 8×8×1 k-mesh and 4000 k-points in the Brillion zone. The TE properties, such as, Seebeck coefficient, electrical conductivity, and thermal conductivity was calculated using Boltzmann transport theory based on BoltzTraP [4, 6], the equation can be written as;

$$\sigma(T;\mu) = \frac{1}{\Omega} \int \sigma(\varepsilon) \left| -\frac{\partial f_{\mu}(T;\varepsilon)}{\partial \varepsilon} \right| d\varepsilon , \qquad (1)$$

written as;  

$$\sigma(T;\mu) = \frac{1}{\Omega} \int \sigma(\varepsilon) \left[ -\frac{\partial f_{\mu}(T;\varepsilon)}{\partial \varepsilon} \right] d\varepsilon , \qquad (1)$$

$$v(T;\mu) = \frac{1}{eT\Omega} \int \sigma(\varepsilon)(\varepsilon - \mu) \left[ -\frac{\partial f_{\mu}(T;\varepsilon)}{\partial \varepsilon} \right] d\varepsilon \qquad (2)$$

$$S = E(\nabla T)^{-1} = (\sigma^{-1})v \qquad (3)$$

$$S = E(\nabla T)^{-1} = (\sigma^{-1})\nu$$
 (3)

where  $\Omega$  is unit cell,  $\nu$  is the band velocity,  $\mu$  is the chemical potential,  $\varepsilon$  is band energy,  $f_{\mu}$  is the Fermi function and E is electromotive force. Since they are no experimental data, we only give the electrical conductivity, electronic thermal conductivity and thermoelectric power factor in terms of  $\tau$ . The electronic thermal conductivity has similar outlines to electrical conductivity, which is because that the electrical thermal conductivity is connected with electrical conductivity by the Wiedemann-Franz law;

$$\kappa_{e} = \sigma L T = \frac{\kappa_{e}}{\tau} \times \tau \tag{4}$$

$$K = K_{\varrho} + K_{L} \tag{5}$$

where L is the Lorenz number,  $\kappa$  is total thermal conductivity  $\tau$  is empirical scattering time and  $\tau$ , is true scattering time.

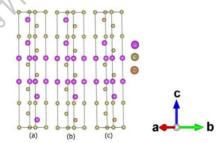


Fig. 1. Cluster atom model (a)  $Bi_2Te_3$ , (b)  $Bi_{1.67}Sb_{0.33}Te_3$ , and (c)  $Bi_{1.33}Sb_{0.67}Te_3$ .

#### 3. Results and discussion

The total density of states of  $Bi_{2-x}Sb_xTe_3$  (0, 0.33, and 0.67) are shows in Fig. 2. The formation energy of  $Bi_{2-x}Sb_xTe_3$  (x=0,0.33,0.67) as shown in Table 1. Both results show semiconductor behavior composed of electron and hole motive, which effected to thermoelectric properties.

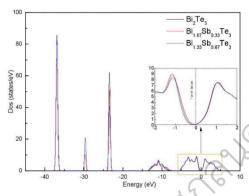


Fig. 2. Total density of states of  $Bi_{2x}Sb_xTe_3$  (x = 0, 0.33, 0.67).

The Seebeck coefficient with respected to scattering time of  $Bi_{2x}Sb_xTe_3$  (x=0,0.33,0.67) at temperature of 300 K as shown in Fig. 3(a). The maximum value of Seebeck coefficient is 536, 436 and 327  $\mu$ V K<sup>-1</sup> occurs at  $\mu=-0.22,-0.26$  and -0.29 eV. The Sb substitutes at Bi has been decreased Seebeck coefficient and showed the p and p-type semiconductors. Electrical conductivity of  $Bi_{2x}Sb_xTe_3$  (x=0,0.33,0.67) was almost independent of temperature as shown in Fig. 3(b). The maximum value of electrical conductivity of  $1.502 \times 10^{20}, 1.35 \times 10^{20}$  and  $1.41 \times 10^{20} \Omega^{-1} m^{-1} s^{-1}$  occurs at  $\mu=-2.4,-2.0$  and -1.9 eV at temperature of 300 K. The entire substitutes Sb at Bi is decrease electrical conductivity from  $1.505 \times 10^{20} \Omega^{-1} m^{-1} s^{-1}$  to  $1.39 \times 10^{20}$  and  $1.41 \times 10^{20}$  to  $1.34 \times 10^{20} \Omega^{-1} m^{-1} s^{-1}$  (at 300 K).

Table 1. The formation energy of  $Bi_{2x}Sb_xTe_3$  (x = 0, 0.33, 0.67).

| Formation energy            | $Bi_2Te_3(Ry)$ | $Bi_{1.67}Sb_{0.33}Te_3$ (Ry) | $Bi_{1.33}Sb_{0.67}Te_3(Ry)$ |
|-----------------------------|----------------|-------------------------------|------------------------------|
| Total energy                | -100048.04     | -9617.47                      | -9186.90                     |
| Total all-electron energy   | -380673.68     | -350575.03                    | -320476.37                   |
| One-electron contribution   | -1281.75       | -1314.63                      | -1347.50                     |
| Hartree contribution        | 771.13         | 785.60                        | 800.06                       |
| XC contribution             | -657.69        | -674.83                       | -691.96                      |
| Ewald contribution          | -1648.60       | -1648.60                      | -1648.60                     |
| One-center paw contribution | -7231.13       | -6765.02                      | -6298.92                     |

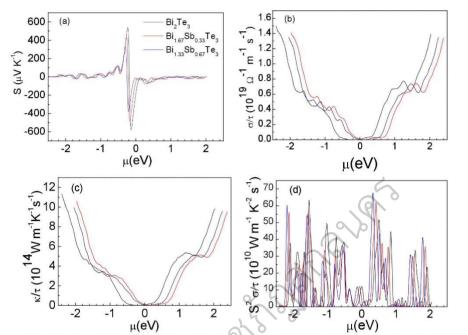


Fig. 3. (a) The (a) Seebeck coefficient, (b) electrical conductivity, (c) thermal conductivity and (d) power factor of  $Bi_{2\alpha}Sb_xTe_3$  (x = 0, 0.33, 0.67) with respected scattering time.

The substitution Sb at Bi as 1 atom and 2 atom are found to value electronic thermal conductivity decrease from  $11.32\times 10^{14}$  to  $9.82\times 10^{14}$  W m  $^{-1}$  K s,  $10.58\times 10^{14}$  to  $9.50\times 10^{14}$  W m  $^{-1}$  K s and  $9.92\times 10^{14}$  to  $9.54\times 10^{14}$  W m  $^{-1}$  K s at (300 K) as shown in Fig. 3(c). The thermoelectric power factor of  $Bi_{2x}Sb_xTe_3$  (x = 0, 0.33, 0.67) plot at temperature 300 K as shown in Fig. 3(d). The maximum value of thermoelectric power factor of (~63.33  $\times 10^{10}$ ,  $61.81\times 10^{10}$  and  $67.752\times 10^{10}$  W m  $^{-1}$  K  $^{-2}$  s  $^{-1}$ ) at  $\mu=-1.54, 0.41, 0.32$  eV.

### 4. Conclusion

The theoretical calculated electronic structure and thermoelectric properties of  $Bi_{2\times}Sb_xTe_3$  (x = 0, 0.33, 0.67) have been investigated by using first principle calculation and Boltzmann transport theory. The thermoelectric properties substitutes Sb at Bi have been increased electrical conductivity and electronic thermal conductivity, while it decreased power factor at 300 K.

## Acknowledgements

This work was supported by Research and Researchers for Industry (RRi) find number MSD59I0075.

#### References

- B.Y. Yavorsky, N.F.Hinsche, I. Mertig, P. Zahn, J. Phys Rev B. 84 (2011) 165208.
   A. Jayaraman, A. B. Kademane, M. Molli, J. Comp Mater Sci. 109 (2015) 34-40.
   H. J. Goldsmid, J. Materials. 7 (2014) 2577-2592.

- [4] K.H.M. Georg, D. J. Singh, J. Comput Phys Commun. 175 (2006) 67–71.
- T.J. Scheidemantel, C. Ambrosch-Draxl, T. Thonhauser, J.V. Badding, J.O. Sofo, J. Phys Rev B. 68 (2003) 125210.
   P. Giannozzi, S. Baroni, N. Bonini, M. Calandra, R. Car, C. Cavazzoni, D. Ceresoli, G. L. Chiarotti, M. Cococcioni, I. Dabo, A. D. Corso, S.
- de Gironcoli, S. Fabris, G. Fratesi, R. Gebauer, U. Gerstmann, C. Gougoussis, A. Kokalj, M. Lazzeri, L. Martin-Samos, N. Marzari, F. Mauri, R. Mazzarello, S. Paolini, A. Pasquarello, L. Paulatto, C. Sbraccia, S. Scandolo, G. Sclauzero, A. P. Seitsonen, A. Smogunov, P. Umari, R. M. Wentzcovitch, J. Physics Condensed Matter. 21 (2009) 395502.
   J. P. Perdew, K. Burke, M. Ernzerhof, J. Phys Rev Lett. 77 (1996) 3865.

## Co- Author

- P. Pilasuta, K. Singsoog, S. Paengsona, W. Namhongsas, S. Ruamruk, L. Sripasudab, S. suwannee, and T. Seetawan. Thermoelectric power generation of p-Ca3Co4O9/n-Zn<sub>0.98</sub>Al<sub>0.02</sub>O lesgs" 4<sup>th</sup> Southeast Asia Confernce on Thermoelectrics 2016. 15.18 December 2016, Sea Garden Hotel, Danang, Vietnam
- 2. Kunchit Singsoog, Panida Pilasuta, Supasit Paengson, Wanachaporn Namhonngsa Weerasak Charoenrat, Surasak Ruamruk, Wiruj Impho, Phanuwat Wongsangnoi Wasana Kasemsin, Tosawan Seetawan, and T. Seetawan THE EFFECT OF SILVER AND BISMUTH DOPED MA2SI ON CRYSTAL TRUCTURE AND THERMOELECTRIC PROPERTIES. Journal of Materials Science and Applied Energy. 6(1), 102–105 (2017) (TCI2)